Holographic Entanglement Entropy For WAdS₃

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Introduction

Guica-Hartman-WS-Strominger Bredberg-Hartman-WS-Strominger Castro-Maloney-Strominger





Introduction

Anninos-Li-Padi-WS-Strominger Compere-Guica-Rodriguez Detournay-Compere Hofman-Strominger Deournay-Hartman-Hofman Compere-WS-Strominger









Review of WAdS3

• EE in WAdS/CFT

• EE in WAdS/WCFT



WAdS₃ holography



WAdS ₃ /WCFT: Y	Vir _L - Kac-Moody _L	Detournay-Compere
Evidence:	$S_{BH} = S_{DHH}$	Deournay-Hartman-Hofman
EE:WS-Wen-Xu	$? = S_{CHI}$	Castro–Hofman–Iqbal



Compere-WS-Strominger



Entanglement entropy in WAdS/CFT

- Assumptions: WAdS/CFT
- Method: Lewkowycz-Maldacena adapted
- Results: The bulk calculation agrees with the CFT expectation



The generalized gravitational entropy

• Replica trick extended to the bulk

$$S_{EE} = -n\partial_n [\log \mathcal{Z}_n - n \log \mathcal{Z}_1]|_{n=1}$$

Lewkowycz-Maldacena

AdS/CFT, Einstein

generalized gravitational entropy

$$= \partial_n (n I[g_n/Z_n] - n I[g_1])|_{n=1} = \partial_n I_{bk}[g_n/Z_n]|_{n=1}$$

• Selecting a special curve

$$ds_{e,n}^2 = n^2 dr^2 + r^2 d\tau^2 + (g_{ij} + 2K_{aij}^{(n)}x^a)dy^i dy^j + \cdots$$

Calculate the generalized gravitational entropy

 $S_{EE} = \int_{\gamma_{\mathcal{A}} \times S^1} \Theta(\phi_i, \partial_n \phi_i) \Big|_{n \to 1, r \to 0}$ presymplectic structure

 $egin{array}{c} n
ightarrow 1, \ \mathcal{K}_a^{(n)} = 0 & ext{Extremal!} \end{array}$

$$S_{EE} = S_{RT}$$

Main assumption: holographic duality exist Role of consistent asymptotic boundary conditions

AdS, Einstein

• set up a dictionary bulk metric bulk metric

$$ds_0^2 = \sigma^{-2} (d\sigma^2 + \gamma_{ij}^{(0)} dx^i dx^j) + h_{\mu\nu} dx^{\mu} dx^{\nu} +$$

• Replica trick extended to the bulk

$$S_{EE} = \partial_n (nI[g_n/Z_n] - nI[g_1])|_{n=1} = \partial_n I_{bk}[g_n/Z_n]|_{n=1}$$

• Selecting a special curve

At n=1, the near cone expansion has to be compatible with the asymptotic expansion

boundary terms only cancels out when the two

metrics satisfy the same asymptotic b.c.

$$ds_{e,n}^2 = n^2 dr^2 + r^2 d\tau^2 + (g_{ij} + 2K_{aij}^{(n)}x^a)dy^i dy^j + \cdots$$

· Calculate the generalized gravitational entropy

 $S_{EE} = \int_{\gamma_{\mathcal{A}} \times S^1} \Theta(\phi_i, \partial_n \phi_i) \Big|_{n \to 1, r \to 0}$

Ambiguity in the presymplectic structure



The generalized gravitational entropy in

WS-Wen-Xu

Consistent boundary conditions

$$ds^{2} = ds_{0}^{2} + warping, \qquad boundary metric$$

$$ds_{0}^{2} = \sigma^{-2}(d\sigma^{2} + \gamma_{ij}^{(0)}dx^{i}dx^{j}) + h_{\mu\nu}dx^{\mu}dx^{\nu}\cdots$$

Selecting a special curve $ds^2 = ds^2 + tilting$

$$ds^{-} = ds^{-}_{e} + tilting,$$

$$ds^{2}_{e} = dr^{2} + r^{2}d\tau^{2} + (\tilde{g}_{ij} + 2K_{aij}x^{a})dy^{i}dy^{j}\cdots$$

Compere-Guica-Rodriguiz

$$n
ightarrow 1$$
, $K_a^{(n)}=0$

NOT geodesic in WAdS, But geodesic in AdS

Compatibility
$$ds_0^2|_{r\to 0} = ds_e^2|_{\sigma\to 0}$$
,
warping $|_{r\to 0} = tilting|_{\sigma\to 0}$,
 $S_{EE} = \frac{\tilde{L}}{4G} = S_{q,m} \neq S_{RT}$



Main Modifications to Lewkowycz-Maldacena for WAdS/CFT

- the expansion near the special surface has to be compatible with the asymptotic expansion;
- periodic conditions are imposed to coordinates in the phase space with diagonalized symplectic structure, not to all fields appearing in the action;
- evaluating the entanglement functional using the boundary term method amounts to evaluating the presymplectic structure at the special surface, where some additional exact form may contribute.



Entanglement entropy in WAdS/WCFT

- Assumptions: WAdS/WCFT, AdS/WCFT
- Method: Casini-Huerta-Myers approach adapted
- Results: The bulk calculation agrees with the WCFT calculation given by Castro-Hofman-Iqbal



CFT $z \rightarrow f(z), \ \overline{z} \rightarrow g(\overline{z})$ $x \to f(x), \qquad t \to t + g(x)$ WCFT

Detournay-Compere Hofman-Strominger Deournay-Hartman-Hofman

By carefully analyzing the moduli properties and keeping track of the anomalis, a Cardy-like formula was derived

moduli transformation

$$S_{WCFT} = -\frac{4\pi i P_0 P_0^{vac}}{k} + 4\pi \sqrt{-\left(L_0^{vac} - \frac{(P_0^{vac})^2}{k}\right)\left(L_0 - \frac{P_0^2}{k}\right)}$$



EE on WCFT

Casini-Huerta-Myers Castro-Hofman-Iqbal

• EE on an interval

$$\mathcal{D}:$$
 $(T,X) \in \left[(-\frac{l_T}{2}, -\frac{l_X}{2}), (\frac{l_T}{2}, \frac{l_X}{2}) \right].$

warped conformal mapping

$$\frac{\tanh\frac{\pi X}{\beta}}{\tanh\frac{l_X\pi}{2\beta}} = \tanh\frac{\pi x}{\kappa}, \qquad T + \left(\frac{\bar{\beta}}{\beta} - \frac{\alpha}{\beta}\right)X = t + \left(\frac{\bar{\kappa}}{\kappa} - \frac{\alpha}{\kappa}\right)x,$$

• EE on \mathcal{D} = Thermal entropy on \mathcal{H}

$$S_{EE} = iP_0^{vac}(\Delta T + \frac{\bar{\beta} - \pi}{\beta}\Delta X) + (iP_0^{vac} - 4L_0^{vac})\log\left(\frac{\beta}{\pi\epsilon}\sinh\frac{\Delta X\pi}{\beta}\right)$$

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A bulk calculation

- WAdS with coordinates (U,V, ρ)
- Bulk coordinate transformation preserves two explicit U(1) isometries

$$\begin{split} u =& \frac{1}{4} \log \left(\frac{(l_U + 2U)^2 \rho^2 - 1}{(l_U - 2U)^2 \rho^2 - 1} \right) \qquad v =& \frac{1}{4} \log \left(\frac{(1 + 2\rho U)^2 - l_U^2 \rho^2}{(1 - 2\rho U)^2 - l_U^2 \rho^2} \right) + V \,, \\ \text{SL(2,R)x U(1) quotient} \qquad r =& \frac{1 + \rho^2 \left(l_U^2 - 4U^2 \right)}{2 l_U \rho} \,, \end{split}$$

• Black string with coordinates (u,v, r) and thermal entopy:

$$S_{thermal} = \frac{\ell}{4G} \Delta V + \frac{\ell}{4G} \log \frac{\Delta U}{\eta}$$
.





Summary of our results

 Assuming WAdS/CFT, we take the LM approach, and derive a holographic calculation for the entanglement entropy. The bulk result agrees with the CFT results.

 Assuming WAdS/WCFT, we take the CHM approach, and derive a holographic calculation for entanglement entropy. The bulk result agrees with the WCFT results.



Thank you!

